Space-time Grid Connections for Finite Integration Method

Kota Arai, Takeshi Mifune, Member, IEEE, Tetsuji Matsuo, Member, IEEE

Graduate School of Engineering, Kyoto University, Kyoto 615-8510, Japan, araikota@fem.kuee.kyoto-u.ac.jp

The perfectly matched layer (PML) absorbing boundary is implemented for the space-time finite integration (FI) method. The subgrid connection in 3D and 4D space-time is discussed. The computational accuracy of several types of 3D and 4D space-time subgrid methods are evaluated using the PML.

Index Terms- Boundary condition, electromagnetic wave absorption, finite integration method, time-domain analysis.

I. INTRODUCTION

THE ELECTROMAGNETIC field analysis of fine structure of sub-wavelength scale is required for advanced electronic and optical devices [1]. The analysis of these devices using the conventional FDTD method [2] suffers large computational cost because the spatial grid should be refined uniformly unless a sophisticated subgrid method [2] is used. The space-time finite integration (FI) method [3]-[4] achieves efficient electromagnetic field computation using an adaptive time-step. Ref. [5] developed 3D and 4D space-time subgrid method for the adaptive grid construction and compared with the subgrid scheme [6] derived from the spatial FI method [7]. However, the computational accuracy of space-time subgrid scheme has not yet been fully examined because only the periodic spatial boundary condition was implemented.

This paper develops the connection scheme to the perfectly matched layer (PML) [8] for the space-time FI method and discusses the 3D and 4D connections to space-time subgrids.

II. SPACE-TIME FINITE INTEGRATION METHOD

The coordinate system is denoted by $(w, x, y, z) = (x^0, x^1, x^2, x^3)$ where w = ct, $c = 1/\sqrt{\varepsilon_0 \mu_0}$ and ε_0 and μ_0 are respectively the permittivity and permeability of vacuum. The integral forms of Maxwell equations [4] without source terms are

$$\oint_{\partial\Omega_p} F = 0, \oint_{\partial\Omega_d} G = 0, \tag{1}$$

$$F = -\sum_{i=1}^{3} E_i dx^0 dx^i + \sum_{j=1}^{3} cB_j dx^k dx^l , \qquad (2)$$

$$G = \sum_{i=1}^{3} H_i dx^0 dx^i + \sum_{j=1}^{3} cD_j dx^k dx^l , \qquad (3)$$

where (j, k, l) is a cyclic permutation of (1, 2, 3), and Ω_p and Ω_d are hypersurfaces in space-time ; $\partial \Omega_p$ and $\partial \Omega_d$ are respectively constituted by the faces of primal and dual grids. The electromagnetic variables in the FI method are defined as

$$f = \int_{S_p} F \, , g = \int_{S_d} G , \tag{4}$$

where S_p and S_d are the faces of the primal and dual grids that constitute $\partial \Omega_p$ and $\partial \Omega_d$. To express the constitutive equation simply, the Hodge dual grid [4] is introduced as

$$\frac{\int_{Sd} c_r dx^0 dx^j}{\int_{Sp} dx^k dx^l} = -\frac{\int_{Sd} dx^k dx^l}{\int_{Sp} c_r dx^0 dx^j} = \kappa , \qquad (5)$$

where $c_r = 1/\sqrt{\varepsilon_r \mu_r}$; κ is a constant determined for each pair of S_p and S_d ; and ε_r and μ_r are respectively the relative permittivity and permeability. From (4) and (5), it follows that $f = Zg/\kappa$, where $Z = \sqrt{\frac{\mu_r \mu_0}{(\varepsilon_r \varepsilon_0)}}$ is the impedance.

III. PML ABSORBING BOUNDARY CONDITION

In the 3D space, all the components of electric flux density and magnetic flux density are divided into two subcomponents respectively such as

$$D_x = D_{xy} + D_{xz} \quad . \tag{6}$$

Using these subcomponents, the space-time FI method updates the electric flux density in the PML as

$$\begin{split} d_{xy,i+\frac{1}{2},j-\frac{1}{2},k-\frac{1}{2}}^{n+\frac{1}{2}} &= \frac{1 - \frac{\Delta W}{2} \sigma_y \sqrt{\frac{\mu_0}{\varepsilon_0}}}{1 + \frac{\Delta W}{2} \sigma_y \sqrt{\frac{\mu_0}{\varepsilon_0}}} d_{xy,i+\frac{1}{2},j-\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}} \\ &+ \frac{1}{1 + \frac{\Delta W}{2} \sigma_y \sqrt{\frac{\mu_0}{\varepsilon_0}}} (h_{z,i,j,k}^n - h_{z,i,j-1,k}^n), \\ d_{xz,i+\frac{1}{2},j-\frac{1}{2},k-\frac{1}{2}}^{n+\frac{1}{2}} &= \frac{1 - \frac{\Delta W}{2} \sigma_z \sqrt{\frac{\mu_0}{\varepsilon_0}}}{1 + \frac{\Delta W}{2} \sigma_z \sqrt{\frac{\mu_0}{\varepsilon_0}}} d_{xz,i+\frac{1}{2},j-\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}} \\ &- \frac{1}{1 + \frac{\Delta W}{2} \sigma_z \sqrt{\frac{\mu_0}{\varepsilon_0}}} (h_{y,i,j,k}^n - h_{y,i,j,k-1}^n), (7) \end{split}$$

where σ_y and σ_z are the electric conductivities in *y*-direction and *z*-direction respectively, Δw is the time-step, the subscripts are spatial indexes for *x*, *y* and *z* directions, and the superscript is the temporal index. The time-marching scheme for the other variables is given similarly.

IV. SPACE-TIME SUBGRID CONNECTION

Ref. [5] proposed straight-type and staircase-type subgrids in the 3D and 4D space-time. However, the subgrid connection in 4D space-time was not discussed in detail.

A. 3D space-time subgrid

The 3D straight-type and staircase-type space-time subgrids are examined with the PML boundary condition. Fig. 1 illustrates the computational domain, where Δx and Δw are set to 1 and 0.5 by normalization and Δl is a free parameter. The normalized initial conditions are $E_x = E_y = 0$ and $B_z =$ $\exp[-(x^2 + y^2)/25]$. Fig. 2 depicts the distributions of discrepancy ΔB_z between B_z obtained employing FDTD method and that obtained using the 3D staircase-type spacetime subgrid at ct = 100. Unphysical wave reflection caused by the subgrid connection is reduced by the optimization of Δl .

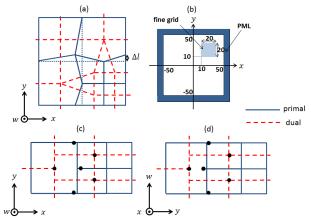


Fig. 1. 3D Space-time subgrid: (a) the straight type subgrid with a free parameter Δl , (b) computational domain, and (c) spatial and (d) space-time subgrid connection of staircase type.

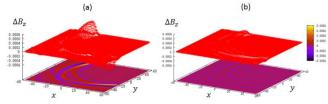


Fig. 2. Discrepancy of B_z compared with FDTD method: (a) $\Delta l = 0.01$, and (b) $\Delta l = 0.1$.

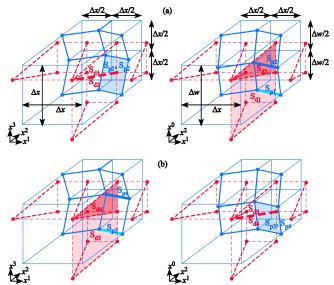


Fig. 3. 4D Straight type subgrid connection: (a) faces S_{p1} , S_{p2} , S_{d1} , S_{d2} , and (b) faces Sp3, Sp4, Sd3, Sd4.

B. 4D Strait type subgrid

Fig. 3 illustrates the straight-type subgrid connection. The faces S_{p1} , S_{p2} and their dual faces S_{d1} , S_{d2} are given as

$$S_{p1} = S_{p2} = \Delta x (dx_2/2 - dx_1/6) \wedge \Delta x (dx_3/2 + dx_1/6)$$
(8)
$$S_{11} = 3\Delta w dx_2/4 \wedge \Delta x (3dx_1/4 + dx_2/4 - dx_2/4)$$
(8)

$$S_{11} = \Delta w dr_0 / 4 - \Delta x (3 dr_1 / 4 + dr_2 / 4 - dr_3 / 4),$$

$$S_{12} = \Delta w dr_0 / 4 \wedge \Delta x (3 dr_1 / 4 + dr_2 / 4 - dr_3 / 4),$$

$$S_{d2} = \Delta w dx_0 / 4^{\wedge} \Delta x (3 dx_1 / 4 + dx_2 / 4 - dx_3 / 4)$$
(9)

The variables $f_i = \int_{\text{Sp}i} F(i = 1, 2)$ mean magnetic fluxes and $g_i =$ $\int_{\text{Sd}i} G(i=1, 2)$ mean magnetomotive forces such as:

$$f_1 = (\Delta x)^2 (3cB_1 + cB_2 - cB_3) / 12$$
(10)

$$g_1 = 3\Delta x \Delta w \left(3H_1 + H_2 - H_3 \right) / 16 . \tag{11}$$

The faces S_{p3} , S_{p4} and their dual faces S_{d3} , S_{d4} are given as

$$_{3} = [\Delta w dx_{0}/2 - (\Delta w)^{2} dx_{1}/6\Delta x] \wedge \Delta x (dx_{2}/2 - dx_{1}/6) ,$$

$$S_{p4} = [\Delta w dx_0/2 - (\Delta w)^2 dx_1/6\Delta x] \wedge \Delta x (dx_2/2 - dx_1/6)$$
(12)

$$S_{d3} = [\Delta x (3dx_1/4 + dx_2/4) - \Delta w dx_0/4]^3 \Delta x dx_3/4,$$

$$S_{d4} = [\Delta x (3dx_1/4 + dx_2/4) - \Delta w dx_0/4]^{\Delta x dx_3/4}.$$
(13)

The dominant component of f_3 and f_4 is the electromotive force but f_3 and f_4 also contain magnetic flux such as:

$$f_3 = \Delta w \left[\Delta x \left(3E_2 - E_1 \right) - \Delta w c B_3 \right] / 12 .$$
 (14)

The dominant component of g_3 and g_4 is the electric flux but g_3 and g_4 also contain magnetomotive force such as:

$$g_3 = 3\Delta x \left[\Delta x (-3cD_2 + cD_1) - \Delta w H_3 \right] / 16 .$$
 (15)

C. 4D Staircase type subgrid

Fig. 4 illustrates the staircase-type subgrid connection. The faces S_{pi} and their dual faces S_{di} (i = 1, ..., 4) are redefined in the staircase type subgrid as

- $S_{p1} = (\Delta x)^2 (dx_2 dx_3/4 + dx_3 dx_1/8 dx_1 dx_2/8),$ $S_{p2} = (\Delta x)^2 (dx_2 dx_3/4 + dz dx_0/24 - dx_1 dx_2/24)$ (16) $S_{d1} = \Delta x \Delta w \left(\frac{3dx_0 dx_1}{8} + \frac{3dx_0 dx_2}{16} - \frac{3dx_0 dx_3}{16} \right),$
- $S_{d2} = \Delta x \Delta w \left(\frac{3dx_0 dx_1}{8} + \frac{dx_0 dx_2}{16} \frac{dx_0 dx_3}{16} \right)$ (17)
- $S_{p3} = \Delta x \Delta w \left(\frac{dx_0 dx_2}{4} \frac{dx_0 dx_1}{8} \right) (\Delta w)^2 \frac{dx_1 dx_2}{8},$
- $S_{p4} = \Delta x \Delta w \left(\frac{dx_0 dx_2}{4} \frac{dx_0 dx_1}{24} \frac{(\Delta w)^2 dx_1 dx_2}{24} \right)$ (18)

$$S_{d3} = (\Delta x)^2 (-3dx_3dx_1/8 + 3dx_2dx_3/16) - 3\Delta x \Delta w dx_0 dx_3/16$$

 $S_{d4} = (\Delta x)^2 (-3dx_3dx_1/8 + dx_2dx_3/16) - \Delta x \Delta w dx_0 dx_3/16 . (19)$

The computational accuracy given by the two types of subgrid connections will be compared in the full paper.

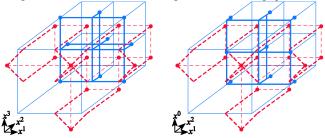


Fig. 4. 4D Staircase type subgrid connection.

V.REFERENCES

- [1] Y. Hao and R. Mittra, "FDTD modeling of metamaterials: Theory and applications," Artech House, 2009.
- A. Taflove and S. C. Hagness, "Computational electrodynamics: The [2] finite difference time-domain method," Artech House, 2005.
- [3] T. Matsuo, "Electromagnetic field computation using space-time grid and finite integration method," IEEE Tran. Magn., vol. 46, pp. 3241-3244 2010
- T. Matsuo, "Space-time finite integration method for electromagnetic [4] field computation," IEEE Trans. Magn., vol. 47, pp. 1530-1533, 2011.
- T. Matsuo, T. Shimoi, J. Kawahara, and T. Mifune, "A simple subgrid [5] scheme using space-time finite integration method," IEEE Trans. Magn., vol. 51, 7201904, 2015.
- P. Thoma and T. Weiland, "A consistent subgridding scheme for the [6] finite difference time domain method," Int. J. Numer. Model., vol. 9, pp. 359-374, 1996.
- T. Weiland, "Time domain electromagnetic field computation with [7] finite difference methods," Int. J. Numer. Model., vol.9, pp.295-319, 1996.
- J.-P. Berenger, "A perfectly matched layer for the absorption of electro-[8] magnetic waves," Journal of Computational Physics, vol. 114, pp. 185-200, 1994.